Chapter 5

Water mass formation, subduction, and the oceanic heat budget

In the first four chapters we developed the concept of Ekman pumping, Rossby wave propagation, and the Sverdrup circulation as the steady-state balance between these two processes. We then showed how the depth-integrated flow in most ocean regions is well described by this concept. However, this is evidently not the full story of the oceanic circulation. The oceans carry heat from the tropics to polar latitudes, and they carry cold water from the poles towards the equator. The details of these transport processes, which are restricted to certain depth ranges, are hidden in the Sverdrup circulation. They have to be resolved if we want to understand the ocean's role in climate variability and climate change.

One thing we do know is that the ocean carries about as much heat towards the poles as the atmosphere does; but since its time scales are so much larger the ocean has a large capacity to act as a damping mechanism for rapid fluctuations in our climate. Conversely, much of the long-term variability of the climate may be related to the ocean as it slowly releases heat stored from earlier rapid climate changes. Topographic details play a large part in determining patterns of ocean heat transport. This is the point of interaction between regional oceanography and geophysical fluid dynamics. It is essential for theoretical studies to know to what degree regional oceanic features have to be part of the modelling process.

Generally speaking, most of the heat and salt exchange with the atmosphere occurs at the ocean surface in a layer which during most of the year is less than 150 m deep. Once a water parcel is removed from the surface layer its temperature and salinity do not change until it rises back up to the surface again, usually many years later. Evidence for this can be found in the temperature and salinity maps of Figures 2.5c and 2.5f for 2000 m depth. The temperature everywhere is less than 4.2°C, even in the tropics; such cold waters can only have acquired their low temperatures in the polar regions and moved into the tropics without much mixing with the warmer waters above. Since independent evidence obtained, for example, from 14C dating sets the residence times of these deeper waters at hundreds of years, this lack of mixing is remarkable. Yet mixing is not entirely absent: Note that the highest temperatures at 2000 m occur just west of the Strait of Gibraltar, where they coincide with very high salinities. The pattern suggests that a "water mass" of relatively warm, salty water enters the Atlantic Ocean from the Mediterranean Sea and moves westward across the ocean, mixing as it goes.

This example illustrates a common feature found throughout the ocean. Water masses with well-defined temperature and salinity characteristics are created by surface processes in specific locations, which then sink and mix slowly with other water masses as they move along. Since these movements are so slow, it is usually unrealistic (and because of the presence of eddy motion impractical in any case) to measure them directly. It is easier to deduce these movements and the strength of mixing from the distribution of the water properties themselves.

The analysis of water mass movements and mixing can assist in understanding the deep oceanic circulation where the estimation of steric height relative to an assumed depth of no motion finds its limits. Steric height estimation provides valuable insights into the workings of the vigorously moving top kilometer of the ocean where accurate knowledge of the depth of no motion is not critical. It cannot resolve the slow flows at great depth. The best hope to achieve a complete description of the circulation in the ocean at all depths is to
combine water mass analysis with our knowledge of the constraints on possible flows posed by geostrophy and Ekman dynamics. This approach, which infers water movement from the effect it has on the distribution of oceanic parameters, is known as "inverse modelling". This is a field of active research, mathematically quite complex and beyond the scope of this book. An understanding of the processes involved in water mass formation and of the life history of water masses is, however, required in regional oceanography.

**Water masses, water types, and T-S diagrams**

We begin with a brief review of concepts and definitions. In the past, oceanographers have used terms such as water mass and water type rather loosely to describe waters with common or outstanding properties. For the purpose of a quantitative description of the transport of water properties in the ocean it is necessary to introduce unambiguous definitions, even though they may not always correspond with earlier uses of the terms. We define a *water mass* as a body of water with a common formation history. An example of water mass formation is the cooling of surface water near the Antarctic continent, particularly in the Weddell Sea, which increases the density and causes the water to sink to great depth. All water which originates from this process shares the same formation history and is called Antarctic Bottom Water. It is found in all oceans well beyond its formation region, extending even into the northern hemisphere.

![Fig. 5.1. Mean T-S diagram and standard deviation ∆S of salinity (for given temperatures) in the eastern Coral Sea, in comparison to water mass definitions in the south Pacific Ocean. Large dots and the heavy line indicate water mass properties in the formation regions, which for all but Surface Water are located far outside the Coral Sea. The standard deviation was determined by comparing stations in the region with a space average and does not include variability in time. Similar standard deviations can be derived for temperature and other properties. Based on Tomczak and Hao (1989).](image)
Common names of known water masses usually relate to their major area of residence. Unfortunately, this can give rise to ambiguity since the same name may be used for a well-defined water mass or simply for water found in a certain region. To avoid this confusion we adopt the convention that water masses are always identified by capitals. For example, "Bottom Water" can stand for Antarctic, Arctic, or other Bottom Water but always refers to a water mass, while water found at the bottom of an oceanic region may be referred to as "bottom water" without implying that it is a known and well-defined water mass. Likewise, we use the term "intermediate water" occasionally for intrusions of water at intermediate depth; in contrast, "Intermediate Water" is used to indicate well-defined water masses.

It is important to note that exclusive occupation of an oceanic region by a single water mass occurs only in the formation regions. As the water masses spread across the ocean they mix, and several water masses are usually present at an oceanic location. However, water masses occupy a measurable volume, which is the sum of the volumes occupied by all its elements regardless of their present whereabouts. It is possible to determine the percentage contribution of all water masses to a given water sample, because the water mass elements retain their properties, in particular their potential temperature and their salinity, when leaving the formation region. Water masses can therefore be identified by plotting temperature against salinity in a so-called T-S diagram. An elementary description of T-S diagrams and their use can be found, for example, in Dietrich et al. (1980). Figure 5.1 is an example of a T-S diagram from a tropical ocean region and shows how observational data can be used to identify water masses from their T-S combinations. It is seen that the properties of Central Water in the Coral Sea correspond closely to those in its formation region, indicating that little mixing with other water masses occurred along its way. In contrast, the intermediate and deep water masses are not present with their original T-S values; their properties are modified by mixing with water above and below, and their presence is indicated by salinity or temperature extrema in the vicinity of the T-S combinations found in their formation regions.

From Figure 5.1 it is seen that T-S relationships alone are insufficient to describe a water mass. Particularly in the upper ocean water masses undergo property changes in response to atmospheric conditions, as indicated in Figure 5.1 by the increase in the standard deviation as the surface is approached. For a full description of a water mass it is necessary to include information about the degree of spatial and long-term variability during its formation, as expressed through its standard deviation. Some water masses, such as the Intermediate Water in Figure 5.1, require only a single T-S combination (T-S point) and a standard deviation; others, such as the Central Water, require a set of T-S combinations, or a T-S relationship, together with a standard deviation envelope. Since data sets to determine standard deviations require observations over several years and are not always available, most books identify water masses by one or more T-S points without standard deviations. Points in the T-S diagram are called water types, and water mass definition points are known as source water types. In reality, very little - if any - of the water belonging to a water mass has exactly the T-S properties of the corresponding source water types. But most T-S values are very close, within the (often unknown) standard deviation. Numerical T-S values for water masses given in later chapters have to be understood in this way. In general, the standard deviation is very small for abyssal water masses but increases rapidly as the upper layers are approached.
The seasonal, tropical, and permanent thermoclines

Most water masses are formed at the ocean surface. This is a region of strong mixing, which produces uniformity of properties above a layer of rapid property change. The term thermocline was occasionally used for this layer in the last two chapters. In the context of water mass formation it is necessary to sharpen our definition of a thermocline.

Oceanographers refer to the surface layer with uniform hydrographic properties as the *surface mixed layer*. This layer is an essential element of heat and freshwater transfer between the atmosphere and the ocean. It usually occupies the uppermost 50 - 150 m or so but can reach much deeper in winter when cooling at the sea surface produces convective overturning of water, releasing heat stored in the ocean to the atmosphere. During spring and summer the mixed layer absorbs heat (moderating the earth's seasonal temperature extremes by storing heat until the following autumn and winter; this aspect is discussed in detail in Chapter 18), and the deep mixed layer from the previous winter is covered by a shallow layer of warm, light water. During this time mixing does not reach very deep, being achieved only by the action of wind waves. Below the layer of active mixing is a zone of rapid transition, where (in most situations) temperature decreases rapidly with depth. This transition layer is called the *seasonal thermocline*. Being the bottom of the surface mixed layer, it is shallow in spring and summer, deep in autumn, and disappears in winter (An example can be seen in Figure 16.27). In the tropics, winter cooling is not strong enough to destroy the seasonal thermocline, and a shallow feature sometimes called the *tropical thermocline* is maintained throughout the year. Figure 5.2 shows the thickness of the surface isothermal layer. It was obtained using the data from Levitus (1982) by extracting the depth where the temperature differed from the temperature at the surface by more than 0.5°C and is representative of the depth of the seasonal thermocline.

Mixed layer dynamics are quite complex, and we refer the reader to other text books (for example Pickard and Emery, 1990) for details. The point of interest here is that turbulent energy levels in the winter mixed layer drop drastically, by a factor of 1000 or more, after it is covered by lighter water during spring - mixing then becomes so slight that the water characteristics established before the turbulence falls off become "frozen". These layers of "fossilized mixing", which retain their signatures for long time spans, are the source layers of most water masses, and their generation is the essence of water mass formation.

The depth range from below the seasonal thermocline to about 1000 m is known as the *permanent or oceanic thermocline*. It is the transition zone from the warm waters of the surface layer to the cold waters of great oceanic depth and provided the model for the interface of our 1½ layer representation of the ocean. The temperature at the upper limit of the permanent thermocline depends on latitude, reaching from well above 20°C in the tropics to just above 15°C in temperate regions; at the lower limit temperatures are rather uniform around 4 - 6°C depending on the particular ocean. Wherever the word thermocline is used without further specification in this book, the term refers to the permanent thermocline. Again, a detailed description of the various thermoclines and their seasonal life cycle can be found in other textbooks (Dietrich et al., 1980; Pickard and Emery, 1990).
Fig. 5.2. Mean depth of the surface isothermal layer (m). (a) August - October, (b) February - April. Contouring levels are 10 m (dashed line), 25 m, 50 m, 75 m, 100 m, 250 m. Adapted from Sprintall and Tomczak (1990).

Subduction

What maintains the permanent thermocline and prevents its erosion from mixing with the waters below and above? The principal factors are the combination of water mass forma-
Figure 4.3 tells us that the subtropics are a region of negative curl($\tau/f$), which means that water is pumped downwards. As this water is not denser than the underlying water, it gets injected into intermediate depths, following the isopycnal surface of its own density. This process, which is known as subduction and illustrated in Figure 5.3, is responsible for the formation of the water masses in the permanent thermocline. Its intensity varies with the seasons, partly in response to variations in the strength of the Ekman pumping but mainly because of the seasonal development of the seasonal thermocline: The summer mixed layer depth is shallower than the depth of the winter mixed layer; the water trapped between (the fossilized mixing region) is available for subduction. If, for example, in late autumn and winter the bottom of the mixed layer ($z_1$ in Figure 5.3) progresses downward faster than the surface water moves downward as a result of Ekman pumping, water pumped from the surface during summer is caught by the expanding surface mixed layer before it can escape into the permanent thermocline, and the properties of the water subducted into the permanent thermocline are determined by the surface water properties during late winter only. In other words, although subduction is a permanent process, water mass formation occurs only in late autumn and winter. This can be verified by comparing the properties of the surface mixed layer in the Subtropical Convergence with those of the permanent thermocline in the tropics, i.e. by comparing T-S diagrams along the lines $ABCD$ and $A'B'C'D'$ in Figure 5.3. An example of such a comparison from the Indian Ocean is shown in Figure 5.4; it demonstrates that over the T-S range of the thermocline the two T-S diagrams are nearly identical in late winter (August - October) but differ during all other seasons.

![Fig. 5.3. Sketch of water mass formation by subduction in the Subtropical Convergence. The T-S diagram shows both the meridional variation of temperature and salinity between stations $A$ and $D$, and the vertical variation equatorward of station $D$ from the surface down along the line $A'B'C'D'$. For more detail, see text.](image)

Water masses subducted into the thermocline are commonly known as Central Water. The term was introduced sixty years ago to differentiate between thermocline water of the
central north Atlantic Ocean (now known as North Atlantic Central Water) and water from the shelf area to the west. It is now used to identify thermocline water masses in all three oceans.

Studies of T-S diagrams and of property distributions on isopycnal surfaces led to the conclusion that mixing across isopycnal surfaces is generally much weaker than mixing on isopycnal surfaces. The principle does not apply universally; western boundary currents, the Equatorial Undercurrent, and frontal regions are among the regions where mixing across density surfaces, or diapycnal mixing, contributes significantly to the exchange of properties. But in general, and over vast ocean regions, it is safe to neglect diapycnal mixing as a first guess for the oceanic circulation.

Figure 5.5 summarizes the discussion of thermoclines and subduction. Intermediate Water, which spreads just below the permanent thermocline, is also produced by subduction. Although the driving agent is not Ekman pumping but mixing and convection in the region between two strong currents the mechanism is the same, movement along isopycnal surfaces towards the equator.

**The barrier layer**

Traditionally, it has been assumed that the depth over which the temperature is uniform can be used as an indication of the depth of the mixed layer, i.e. the layer affected by surface mixing processes. This assumption probably developed more out of necessity than physical argument, since many upper ocean observations lack information on salinity and some way had to be found to determine the depth of the surface mixed layer from temperature information alone. The better data base of today allows us to check that assumption, thereby gaining a better understanding of the processes of water mass formation. Figure 5.6
shows the thickness of the surface isopycnal layer, i.e. the layer over which density does not change. This map was obtained by determining the depth where density is larger than the density at the surface by an amount which corresponds to the temperature change of 0.5°C used in the construction of Figure 5.2. This amount depends strongly on temperature and also on salinity. To make sure that the distributions of Figures 5.2 and 5.6 are comparable, the density increment is not constant across the map but was evaluated locally from the surface temperature and salinity. Figure 5.7 shows a map of the differences between the calculated isothermal and isopycnal layer thicknesses.

![Diagram](image)

Fig. 5.5. Meridional section through a hypothetical ocean, showing the permanent thermocline, the seasonal and tropical thermoclines, the various surface convergences and divergences, and the major water masses. Note the scale change at 1000 m depth, which underestimates the volume of Deep and Bottom Water. On the other hand, the importance of the tropics does not come out in this graph since it does not show the convergence of meridians towards the poles.

If the classical assumption were correct, the differences shown in Figure 5.7 should be close to zero everywhere. There are indeed large ocean areas which display very small differences. But we also notice regions where the difference is clearly not zero. In the tropics, the difference between isothermal and isopycnal layer depth is often positive, indicating a density change within the isothermal layer. The density change is caused by salinity stratification. In these regions, the halocline is the true indicator of mixed layer depth. The tropical surface water has low salinities and high temperatures and therefore very low densities; it spreads in a thin layer over the top of the water column. Subtropical surface water, on the other hand, has high salinities; when this water is subducted in regions of Ekman pumping it cannot be pushed very deep before it spreads sideways, forming a salinity maximum below the surface layer. If the temperature gradient is insignificant, the result of both processes is a halocline within the isothermal layer.
The layer between the halocline and the thermocline is now referred to as the barrier layer, because of its effect on the mixed layer heat budget. The mixed layer receives large amounts of heat from solar radiation. In a steady state situation this source of heat has to be balanced by one or more sinks. In the absence of a barrier layer, i.e. where the surface mixed layer...
reaches down into the thermocline, an important sink is located at the bottom of the mixed layer where cold water is entrained from below. The presence of a barrier layer means that the water entrained into the mixed layer has the same temperature as the water above. There

Fig. 5.7. Mean depth difference (m) between isothermal and isopycnal layer (barrier layer thickness). (a) August - October, (b) February - April. Broken contours indicate negative differences. Adapted from Sprintall and Tomczak (1990).
is therefore no heat flux through the bottom of the mixed layer, and other sinks have to come into play to prevent a permanent rise of mixed layer temperature. The dynamics of the three regions which show a barrier layer through most of the year, the western Pacific, the equatorial Atlantic, and the Bay of Bengal in the Indian Ocean, differ distinctly from each other; they will be discussed in more detail in the chapters which deal with the individual oceans.

Another region of positive difference in isothermal vs isohaline layer depth is in the polar latitudes where a thermocline does not exist: The surface water, constantly cooled by the atmosphere, is so dense that it sinks virtually to the bottom of the ocean (the refinements of this process will be discussed in Chapters 6 and 7), leaving the water column isothermal to great depth. The salinity, on the other hand, is seasonally influenced by freshwater release from the ice shelf and from icebergs. The resulting reduction in density can inhibit the sinking of water, and a halocline can develop. This process, seasonal development of a pycnocline as a result of freshwater flux from the ice regions, is indicated by the large differences in Figure 5.7 south of 50°S and during February to April near Newfoundland and Labrador and south of Alaska and the Aleutian Islands.

The third region of non-zero differences in Figure 5.7 is located in the subtropics and displays negative values, indicating that the first density change below the surface is found at a depth greater than the isothermal layer thickness. This can only occur if the density change produced by the change in temperature across the thermocline is compensated by an appropriate salinity change. To explain this feature we note from Figure 2.5 that in the subtropics where the water of the permanent thermocline is subducted, sea surface temperature and salinity both decrease rapidly towards the poles. The T-S diagram which describes the meridional variation of temperature and salinity across the subduction zone thus nearly follows an isopycnal (compare Figure 5.3). As a consequence, the isothermal layer thickness, say $z_I$ at station $C$, is evaluated correctly by the 0.5°C criterion of Figure 5.2. The equivalent density criterion, however, is not exceeded until the greater depth $z_2$ is reached, because of the compensating salinity effect. It is seen that the difference in layer thickness shown in Figure 5.7 for the subtropics does not, as in the case of the tropics, result from different isothermal and isohaline layer thickness. We might even regard it as an artifact produced by the subduction process. On the other hand, the discussion in Chapters 9, 12, and 15 will show that the negative differences in the subtropics are a reliable indicator for subduction of thermocline water.

The picture that emerges particularly in the subtropical gyres is a set of independent flows on isopycnal surfaces which, when depth-integrated, collectively satisfy the Sverdrup relationship. In the following chapters, we shall use the Sverdrup circulation as a guide for the discussion and fill in the depth-dependence by looking at water mass properties where necessary. One instance where this need will arise is whenever we want to evaluate oceanic transports of heat or salt, because both involve integrals over products of velocity and temperature or salinity which require knowledge of the distribution of currents with depth. To see this, consider the example of warm water moving poleward with a velocity of 0.4 m s$^{-1}$ in the upper 1000 m and cold water moving equatorward with 0.1 m s$^{-1}$ between 1000 m and 5000 m. The net transport of heat is obviously poleward, but the net mass transport is zero. Clearly, the vertically integrated flow patterns of Chapter 4 are inadequate for estimating heat or salt transports.